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## on the lagrange problem of the mean motion of perihelia*

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It is shown that themean motion of perihelia in the Lagrange sense in a non-resonance set is uniformly continuous in the initial phases of the frequency function.

The dynamics of a planetary system such as the solar system is examined. To a first approximation of perturbation theory, when the squares of the orbit eccentricities can be neglected compared with the eccentricities themselves, the evolution of the Laplace vector is described by the function

$$
A(t)=\sum_{m=\pi=0}^{n} a_{m} \exp \left[2 \pi i\left(\lambda_{m} t+\varphi_{0_{m}}\right)\right]
$$

where the constants $a_{m}, \lambda_{m}, \varphi_{0 m}$ are determined in terms of the planet mass and the inftial conditions. The mean motion of the perihelion is defined as

$$
\mu\left(a, h, \varphi_{0}\right)=\lim _{t \rightarrow \infty} \frac{1}{t} \arg A(t)
$$

Lagrange showed that $\mu=2 \pi \lambda_{0}$ if $a_{0} \geqslant a_{1}+a_{2}+\ldots+a_{n+}$ In the mon-trivial case when this condition is not satisfied, the mean motion is calculated for $n=2[1]$ and for arbitrary $n$ $/ 2 /$ for the non-resonance set of frequencies $\lambda=\left(\lambda_{0}, \ldots, \lambda_{n}\right):<k, \lambda>\neq 0, \forall k \in Z^{n+1}, k \neq 0$ and has the form

$$
\begin{aligned}
& \mu(a, \lambda)=2 \pi \sum_{m=0}^{n} \lambda_{m} w^{m}(a) \\
& \sum_{m=0}^{n} w^{m}(a)=1, \quad w^{m}(a) \geqslant 0, \quad n=\left(a_{0}, \ldots, a_{n}\right)
\end{aligned}
$$

The existence of mean motion for an arbitrary "set of frequencies $\lambda$ is proved in $/ 3 /$. Let $a, \lambda$ be certain continuous Eunctions of the parameter $\alpha \equiv Z^{*}$.
 (9 mod 4 )

$$
\lim _{\alpha \rightarrow \alpha_{4}} \mu\left(a(\alpha), \lambda(\alpha), \varphi_{0}\right)=\mu\left(a\left(\alpha_{0}\right), \lambda\left(\alpha_{0}\right)\right)
$$

Remarks. $1^{\circ}$. For fixed $a$ the function $\mu\left(a, \lambda, \varphi_{0}\right)$ is generally discontinuous in $\boldsymbol{o}_{a}$, when $\lambda(\alpha)$ is a resonance vector.
$2^{\circ}$. If the function $A(t)$ vanishes for certain values of the time, its argument is not defined. In such a case it is customary to distinguish the "first" and "left" arguments of the function $A(t)$. When passing through a zero of nultiplicity $p$ the right argument of the function $A(t)$ receives an increment $\pi p$ as $t \rightarrow \infty$, while the left receives the increment ( $-\pi p$ ) /4/. The right and left mean motions $\mu^{+}$and $\mu-/ 3 /$ that are in agreement for non-resonance

[^0]sets of frequencies are defined correspondingly. The assertion formulated is valid for both the left and the right mean motions.

Proof of the assertion. Let $\lambda_{0}(\alpha) \neq 0$. We define $B_{i y} \ldots i_{g}(x) \in T^{n+1} \cap\left\{\varphi_{0}=0\right.$ as the set of points $\Psi_{0}=\left(0, \varphi_{04}, \ldots \varphi_{m}\right)$ such that the corresponding function $A(t)$ satisfies the following conditions:
$1^{\circ}$. A $(t)$ has just $i m$ zeros of multiplicity $m$ in the half-interval $\left[0,1 / \lambda_{0}\right)$.
$2^{\circ}$. Let $\delta^{+}$be the number of zeros $\left\{t_{s}{ }^{+}\right\}$of the function $\operatorname{Im} A(t)$ such that

$$
t_{s}{ }^{+} \in\left[0,1 / \lambda_{0}\right), \quad \operatorname{Im} A\left(t_{s}^{+}+0\right)>0, \quad \operatorname{Im} A\left(t_{s}^{+}-0\right)<0, \quad \text { Re } A\left(t_{s}+\right)<0
$$

$3^{\circ}$. Let $\delta$-be the number of zeros $\left\{t_{s}\right\}$ of the function $\operatorname{Im} A(t)$ such that

$$
t_{s}^{-} \equiv\left(0,1 / \lambda_{0}\right), \operatorname{Im} A\left(t_{s}^{-}+0\right)<0, \operatorname{Im} A\left(t_{s}^{-}-0\right)>0, \operatorname{Re} A\left(t_{s}\right)<0
$$

Then $i=\delta^{-}-\delta^{+}$.
It follows from the Bol'-Weyl construction / $/ 1,2 /$ that for arbitrary initial conditions the mean motion is calculated thus:

$$
\begin{aligned}
& \mu^{ \pm} \pm\left(\alpha, \lambda_{1} \psi_{0}\right)=\lim _{N \rightarrow \infty} \frac{1}{N^{\prime}} \sum_{m=0}^{N-\alpha} \theta^{N}\left(\psi_{0}+m \omega(\alpha), \alpha\right) \\
& \theta \pm(\psi, \alpha)=2 \pi \lambda_{0}(\alpha) \sum_{i, j_{i, \ldots, j_{k}}}\left(i \pm \frac{1}{2} \sum_{m=1}^{g} m i_{m}\right) x\left(\psi, B_{i j, \ldots j_{g}}(\alpha)\right) \\
& \omega(\alpha)=\left(\frac{\lambda_{1}}{\lambda_{0}}, \ldots, \frac{\lambda_{n}}{\lambda_{0}}\right), \psi_{0}=\left(\varphi_{01}-\varphi_{00} \frac{\lambda_{1}}{\lambda_{00}}, \ldots, \varphi_{0_{n}}-\varphi_{00} \frac{\lambda_{n}}{\lambda_{0}}\right)
\end{aligned}
$$

For any $\varepsilon>0$ there exists a neighbourhood $U_{1} E D$ of the point $x_{0}$ and two continuous functions $F_{1}(\psi, \varepsilon), F_{2}(\psi, \varepsilon)$ on $T^{\prime}$ such that

$$
\begin{align*}
& F_{3}(\psi, \varepsilon) \leqslant \theta^{ \pm}(\psi, \alpha) \leqslant F_{i}(\psi, \varepsilon), V(\psi, \alpha) \in T^{n} \times U_{1}  \tag{1}\\
& \int_{T^{n}}\left[F_{2}(\psi, \varepsilon)-F_{1}(\psi, \varepsilon)\right] d \phi \leqslant \varepsilon / 3
\end{align*}
$$

Indeed, if $U_{2}$ is a sufficiently small neighbourhood of $\alpha_{0}$, then the number of zeros of the analytic functions ImA(t) in

$$
\left[0, \operatorname{sign} \lambda_{0} \cdot \sup \left|\frac{1}{\lambda_{0}(\alpha)}\right|\right)
$$

is limited by the constant $\psi$ uniformly in $\alpha \in U_{s}$. The multiplicity of each zero of $A(t)$ is also iimited uniformly in $\alpha \in V_{2}$. Then

$$
\left|\theta^{ \pm}(\psi, \alpha)\right| \leqslant \Gamma, \quad \forall(\psi, \alpha) \equiv 7^{n} \times U_{2}
$$

For fixed $\alpha$ the functions $\theta^{+}(\phi, a)$ and $\theta^{-}(\phi, \alpha)$ agree everywhere except the set

$$
V(\alpha)=U_{i} \sum_{m} U_{m \neq 0} B_{i 3 k+\ldots j g}(\alpha), \quad 0<\operatorname{mes}_{n-1} V(\alpha)<\infty
$$

If $S\left(\alpha, \varepsilon_{1}\right)$ is an $\varepsilon_{1}$-neighbourhood of the set $V(\alpha)$, then for $a \in U_{3}$ where $U_{3}$ is a sufficiently small neighbourhood of $a_{0}$

$$
\begin{aligned}
& V(\alpha) \in S\left(\left(\alpha_{0}, c_{0}\right), \quad \varepsilon_{1}=\frac{\varepsilon}{48 \Gamma \operatorname{mes} s_{n-1} V\left(\alpha_{0}\right)}\right. \\
& \sup _{\alpha=U_{2}} \lambda_{0}(\alpha)-\inf _{\alpha \in i}, \lambda_{0}(\alpha) \leqslant \frac{\varepsilon}{\varepsilon_{i} \pi \gamma} .
\end{aligned}
$$

Let

$$
G^{ \pm}(\psi)= \begin{cases} \pm \Gamma, & \psi \in S\left(\alpha_{0,}, \varepsilon_{1}\right) \\ \sup _{\alpha \in U_{2} \cap V_{1}}\left(\operatorname{ini}_{\alpha \in U_{2} \cap U_{1}}\right)^{\theta^{ \pm}(\psi, \alpha),} & \psi \notin S\left(\alpha_{0}, F_{1}\right)\end{cases}
$$

Then

Now applying the standard method of approximating step functions by continuous functions to $G^{-}$and $G^{*}$, we obtain the functions $F_{1}, F_{\mathrm{s}}$ that satisfy inequalities (1).

The subsequent reasoning is similar to the reasoning in Kozlov's theorem about time means /5/.

According to Weierstrass's theorem on approximation /6/, trigonometric polynomials $Q_{1}(\varphi)$, $Q_{2}(\Psi)$ exist such that

$$
\left|Q_{m}(\psi)-F_{m}(\psi, \varepsilon)\right|<\varepsilon / 6, \quad m=1,2
$$

Then the polynomials $P_{1}(\boldsymbol{\psi})=Q_{1}(\boldsymbol{\psi})-\varepsilon / 6, \quad P_{2}(\boldsymbol{\Psi})=Q_{2}(\boldsymbol{\psi})+\varepsilon 6$ satisfy the conditions

$$
\begin{align*}
& P_{1}(\psi)<\theta^{ \pm}(\psi, \alpha)<P_{2}(\psi), \quad \forall(\psi, \alpha)=T^{n} \times U_{1}  \tag{2}\\
& \int_{T^{n}}\left[P_{2}(\psi)-P_{1}(\psi)\right] d \psi<\varepsilon
\end{align*}
$$

Let $d=\max \left(\operatorname{deg} P_{1}, \operatorname{deg} P_{s}\right)$. Since $\omega\left(\alpha_{0}\right)>\boldsymbol{f} \quad$ for $\alpha=\alpha_{0}<k$ for all $k \in Z^{n}, k \neq 0$, then in a certain neighbourhood $U_{4} \equiv D$ of the point $\alpha_{0}$ the following inequalities hold

$$
\begin{equation*}
\left|\langle\mathbf{k}, \boldsymbol{\omega}(\boldsymbol{\alpha})\rangle-k_{0}\right|>e>0, \quad 0<|\mathbf{k}|=\sum_{m=1}^{n}\left|k_{m}\right| \leqslant d \tag{3}
\end{equation*}
$$

If

$$
g(\downarrow)=\sum_{0 \leqslant|\mathbf{k}| \leqslant d} g_{k} \exp [2 \pi i\langle\mathbf{k}, \phi\rangle]
$$

then upon compliance with the inequalities (3) we have

$$
\lim _{N \rightarrow \infty} \frac{1}{N} \sum_{m=0}^{N-1} g\left(\psi_{0}+m \omega(\alpha)\right)=g_{0}=\int_{T^{n}}^{0} g(\psi) d \psi
$$

Then

$$
\lim _{N \rightarrow \infty} \frac{1}{N} \sum_{m=0}^{N-1} P_{q}\left(\psi_{0}+m \omega(\alpha)\right)=\Lambda_{q}=\int_{T^{n}} P_{q}(\psi) d \psi, q=1,2
$$

By virtue of the first inequality in (2)

$$
\frac{1}{N} \sum_{m=0}^{N-1} P_{1}\left(\psi_{0}+m \omega(\alpha)\right)<\frac{1}{N} \sum_{m=0}^{N-1} \theta^{ \pm}\left(\psi_{0}+m \omega(\alpha), \alpha\right)<\frac{1}{N} \sum_{m=0}^{N-1} P_{2}\left(\psi_{0}+m \omega(\alpha)\right)
$$

and we have in the limit as $\quad N \rightarrow \infty$

$$
\Lambda_{1} \leqslant \mu^{ \pm}\left(\mathrm{a}(\alpha), \lambda(\alpha), \varphi_{0}\right) \leqslant \Lambda_{i}, \quad \alpha \in U_{1} \cap U_{4}=U
$$

But by virtue of the first inequality in (2) for $\alpha=\alpha_{0}$

$$
\Lambda_{1} \leqslant \mu\left(a\left(\alpha_{0}\right), \lambda\left(\alpha_{0}\right)\right) \leqslant \Lambda_{2}
$$

from which by virtue of the second inequality in (2)

$$
\left|\mu^{ \pm}\left(\mathrm{a}(\alpha), \lambda(\alpha), \varphi_{0}\right)-\mu\left(\mathrm{a}\left(\alpha_{0}\right), \lambda\left(\alpha_{0}\right)\right)\right|<\varepsilon, \quad \alpha \in U
$$

which it was required to prove.
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